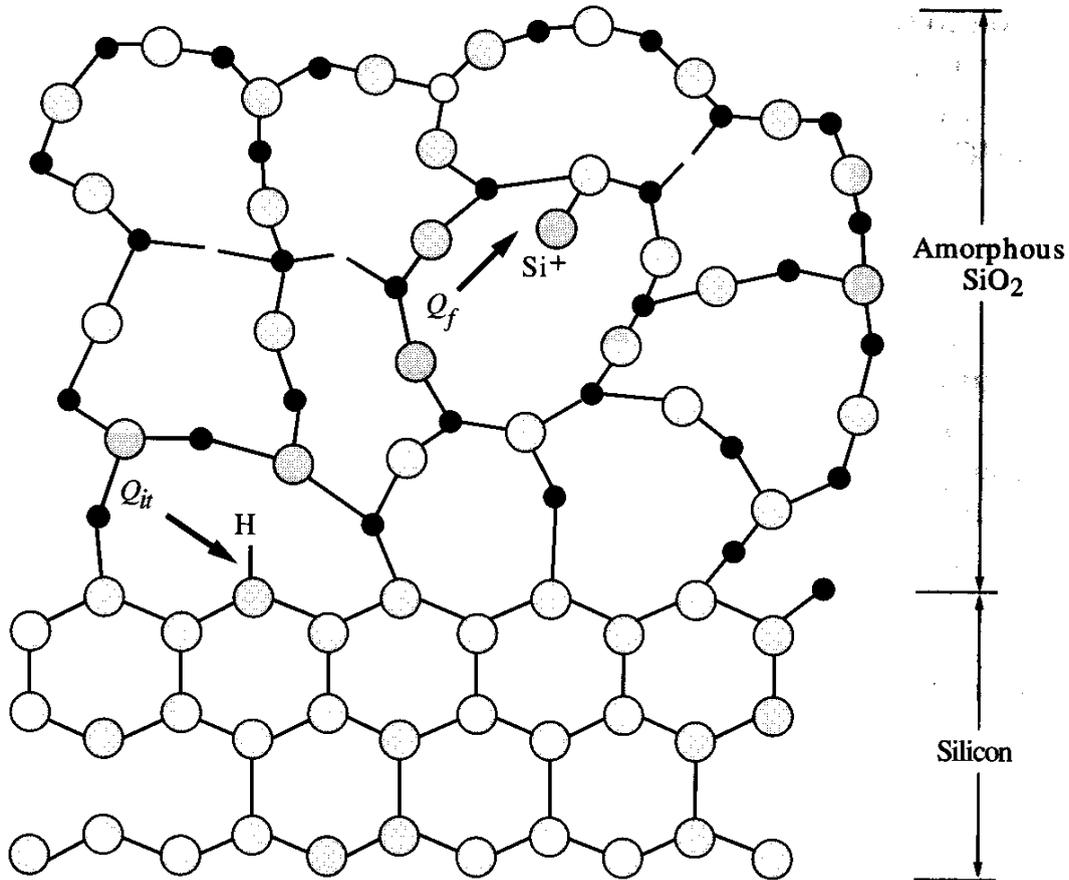
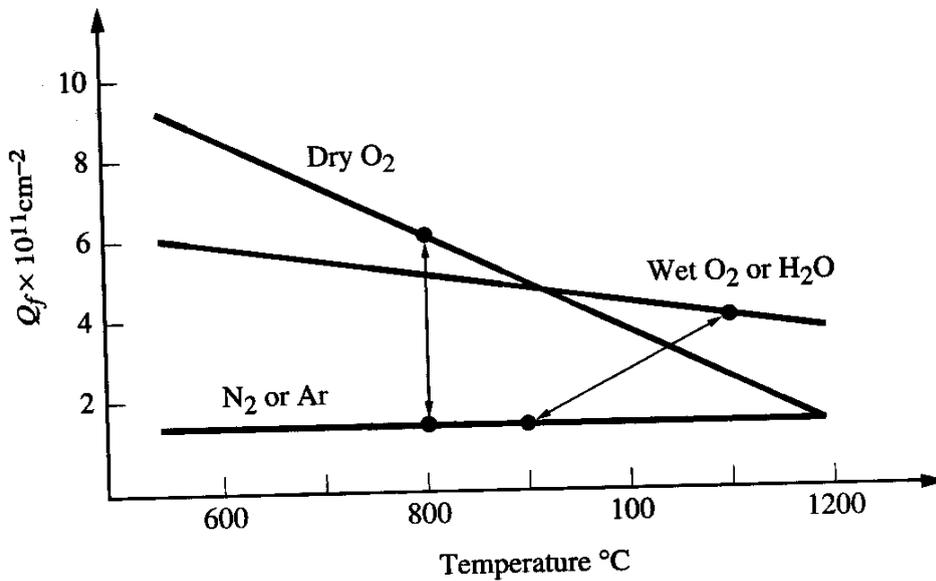


### 6.5.13 Si/SiO<sub>2</sub> Interface charges



- The bulk of the oxide consists mainly of rings of SiO<sub>4</sub> tetrahedra, each with six Si atoms in the ring
  - =Si• corresponding to  $Q_f$  are located physically further away from the interface and thus are unable to trap carriers
- The energy levels corresponding to  $Q_f$  are located outside the bandgap



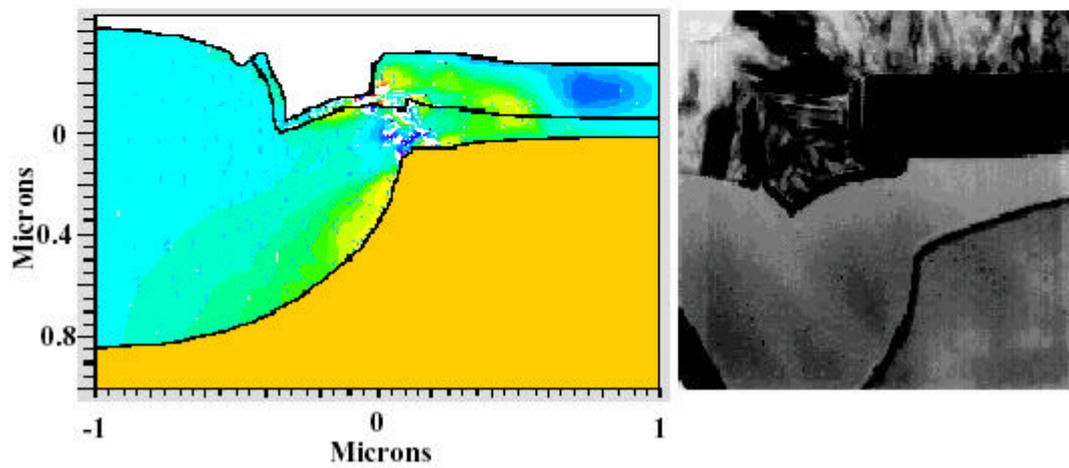
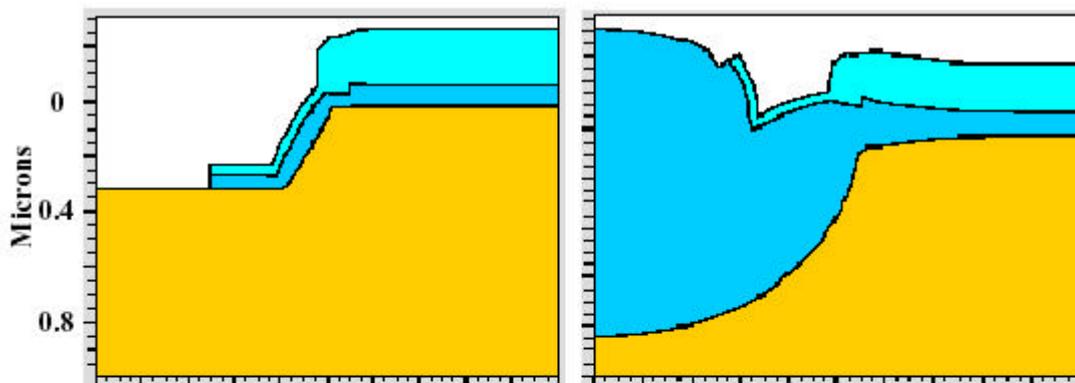
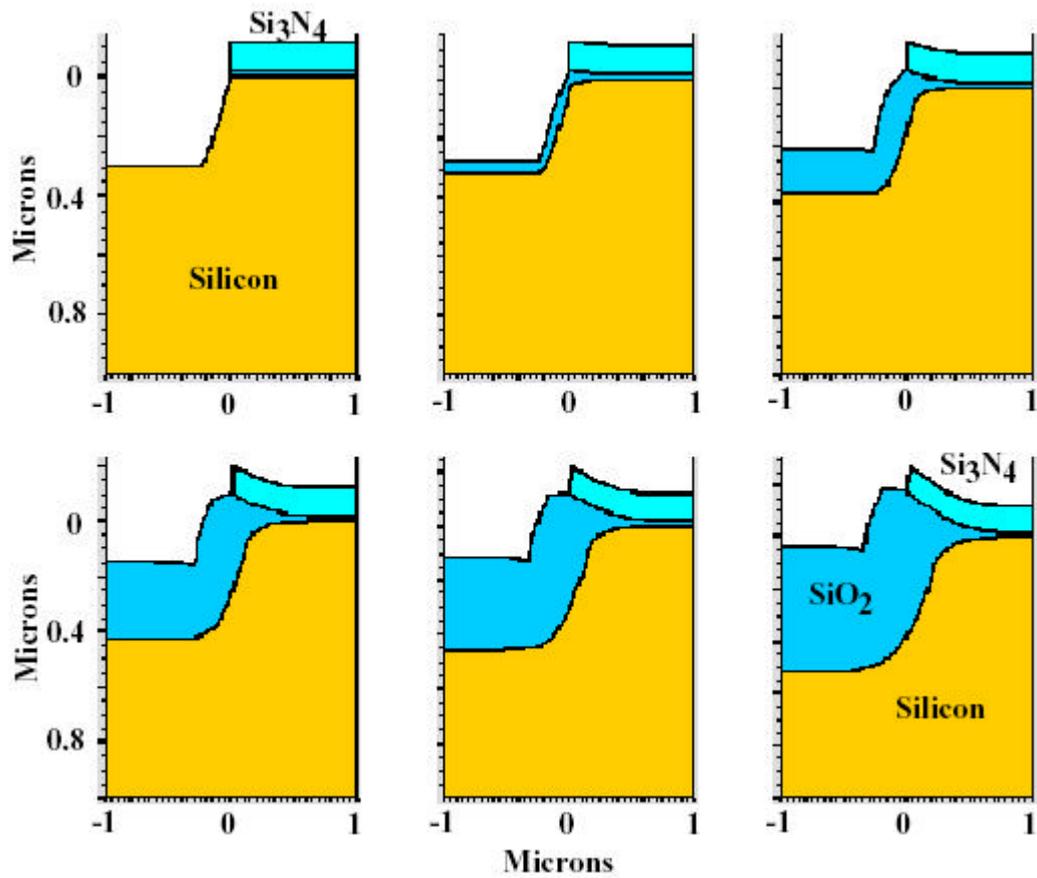
$$Q_f(t) = (Q_f(0) - Q_{fe}) \exp\left(-\frac{t}{\tau}\right) + Q_{fe}$$

- Ar is a better choice than  $\text{N}_2$  for annealing
  - $Q_{fe}$  is essentially independent of temperature
  - $Q_{it}$  states can be passivated through hydrogen annealing at fairly low temperatures whereas  $Q_f$  is unaffected
- $\text{H}_2$  ?  $2\text{H}$  followed by  $=\text{Si}\cdot + \text{H} ? = \text{SiH}$

$$Q_{it} = \frac{Q_{it}(0)}{(1 + kt)^n}$$

- In large area structures such as MOS capacitors, the diffusion of  $\text{H}_2$  can involve significant diffusion times
- Polysilicon gates have the advantage of preventing outdiffusion of  $\text{H}_2$  ? annealing process is more tolerant of high-T exposure

### 6.5.14 Complete Oxidation Module Simulation



## **6.6 Limits and Future Trends in Technologies and Models**

- **Low T**
- **Low T + high P (10-25 Atm)**
- **The ideal properties of the Si/SiO<sub>2</sub> interface can be preserved with deposited oxides either by first growing a thin thermal oxide and then depositing SiO<sub>2</sub> on top of it or by annealing the deposited oxide**
- **“Composite” dielectric layers**
- **Oxynitride dielectrics by exposing thermally grown SiO<sub>2</sub> films to N<sub>2</sub>O or NO ambients, or nitrogen implants and subsequent oxidation**
  - Stronger Si-N bonds ? resistant to hot carrier stressing**
  - Resistant to dopant penetration**
  - Larger K**
- **Change density ~ 10<sup>10</sup>cm<sup>-2</sup> ? one charge / 0.1×0.1μm**



$$J = nqv = nqme = se = \frac{1}{r}e$$

$$r = \frac{e}{J}Wcm$$

In Fig. 7 - 2 (b),

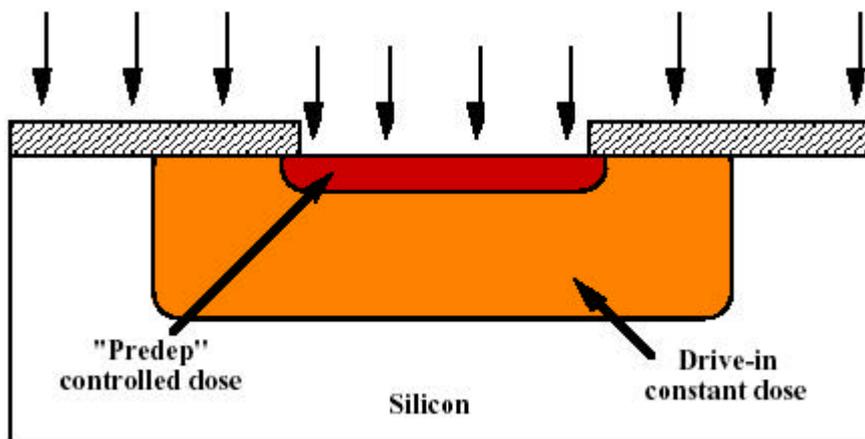
$$R = \frac{r}{x_j} W/\text{square} \circ r_s$$

$$r_s = \frac{1}{Sx_j} = \frac{1}{q \int_0^{x_j} [n(x) - N_B] m [n(x)] dx}$$

- $R_{\text{contact}} + R_{\text{source}} + R_{\text{ext}} < \frac{R_{\text{chan}}}{10}$

Year of 1st DRAM Shipment	1997	1999	2003	2006	2009	2012
Min Feature Size (nm)	250	180	130	100	70	50
DRAM Bits/Chip	256M	1G	4G	16G	64G	256G
Minimum Supply Voltage (volts)	1.8-2.5	1.5-1.8	1.2-1.5	0.9-1.2	0.6-0.9	0.5-0.6
Gate Oxide $T_{\text{ox}}$ Equivalent (nm)	4-5	3-4	2-3	1.5-2	<1.5	<1.0
Sidewall Spacer Thickness $x_w$ (nm)	100-200	72-144	52-104	20-40	7.5-15	5-10
Contact $x_j$ (nm)	100-200	70-140	50-100	40-80	15-30	10-20
$x_j$ at Channel (nm)	50-100	36-72	26-52	20-40	15-30	10-20
Drain Ext Conc ( $\text{cm}^{-3}$ )	$1 \times 10^{18}$	$1 \times 10^{19}$	$1 \times 10^{19}$	$1 \times 10^{20}$	$1 \times 10^{20}$	$1 \times 10^{20}$

## 7.2 Historical Development and Basic Concepts

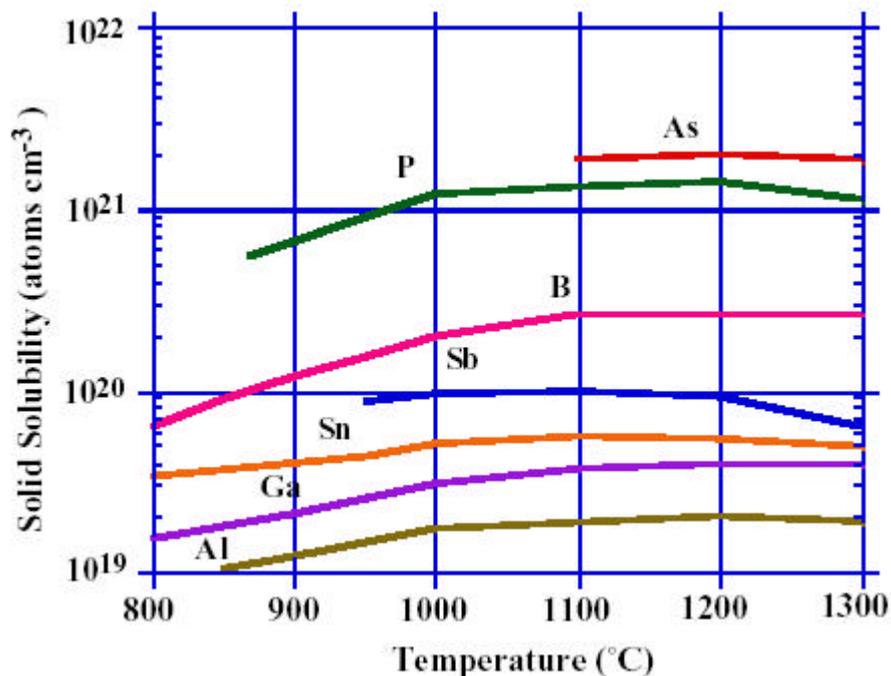


## TED (Transient Enhanced Diffusion)

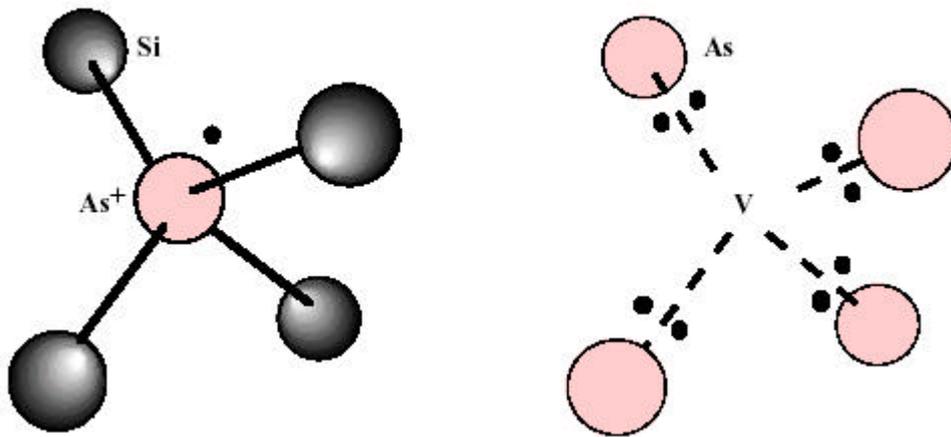
	Ion Implantation and Annealing	Solid/Gas Phase Diffusion
Advantages	Room temperature mask	No damage created by doping
	Precise dose control	Batch fabrication
	$10^{11}$ - $10^{16}$ atoms $\text{cm}^{-2}$ doses	
	Accurate depth control	
Problems	Implant damage enhances diffusion	Usually limited to solid solubility
	Dislocations caused by damage may cause junction leakage	Low surface concentration hard to achieve without a long drive-in
	Implant channeling may affect profile	Low dose predepos very difficult

### 7.2.1 Dopant Solid Solubility

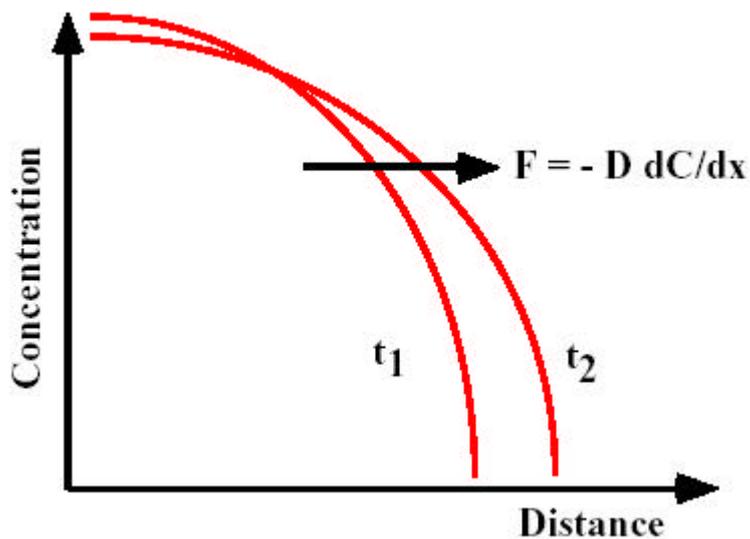
Maximum concentration that can be dissolved in Si under equilibrium conditions, without forming a separate phase



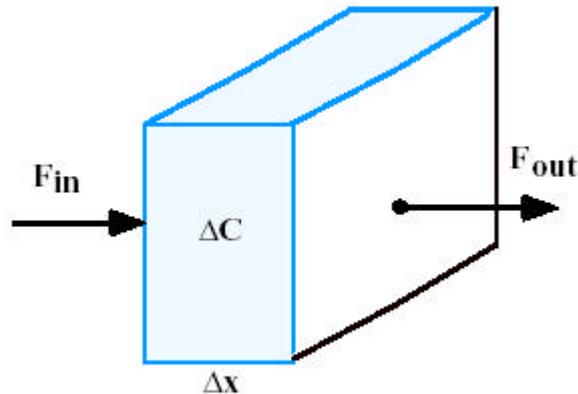
**Electrical solubility limit may be considerably lower than  
The maximum solid solubility ( neutral cluster formation  
With point defects in the Si lattice)**



### 7.2.2 Diffusion from a Macroscopic Viewpoint



- Fick's first law :  $F = -D \frac{\partial C}{\partial x}$
- In a diamond lattice (Si, GaAs) which has cubic symmetry,  $D$  has the same value in all directions



- Fick's second law

$$\frac{DC}{Dt} = \frac{DF}{Dx} = \frac{F_{in} - F_{out}}{Dx}$$

$$\frac{\partial C}{\partial t} = \frac{\partial}{\partial x} \left( D \frac{\partial C}{\partial x} \right)$$

- If  $D$  is a constant

$$\frac{\partial C}{\partial t} = D \frac{\partial^2 C}{\partial x^2}$$

$$\frac{\partial C}{\partial t} = -\tilde{N} \cdot F = \tilde{N} \cdot (D\tilde{N}C)$$

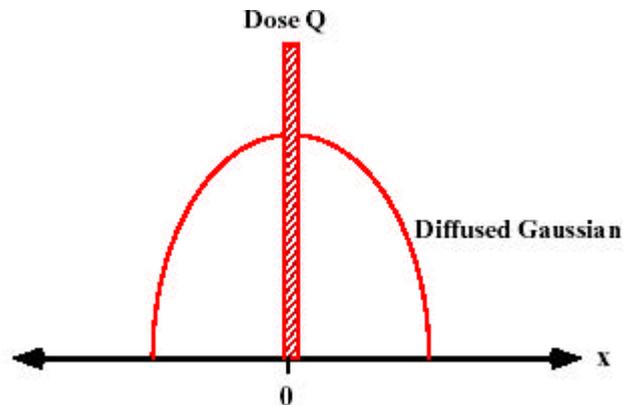
### 7.2.3 Analytic Solutions of the Diffusion Equation

- Steady-state

$$D \frac{\partial^2 C}{\partial x^2} = 0$$

$$C = a + bx$$

## 7.2.4 Gaussian Solution in an Infinite Medium



$C \rightarrow 0$  as  $t \rightarrow 0$  for  $x > 0$

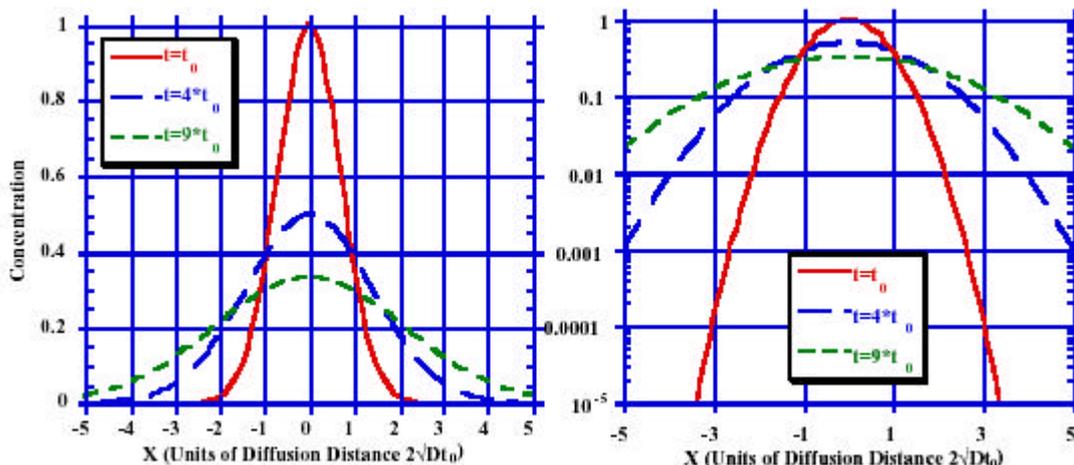
$C \rightarrow \infty$  as  $t \rightarrow 0$  for  $x = 0$

$$\int_0^{\infty} C(x,t) dx = Q$$

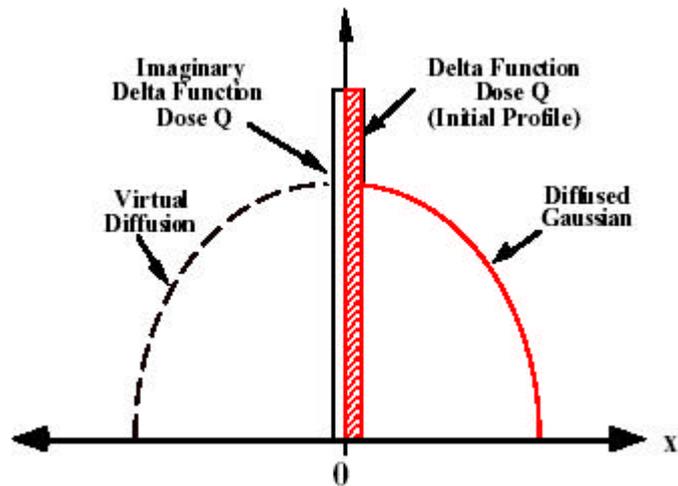
$$C(x,t) = \frac{Q}{2\sqrt{pDt}} \exp\left[-\frac{x^2}{4Dt}\right] = C(0,t) \exp\left[-\frac{x^2}{4Dt}\right]$$

The peak concentration decreases as  $\frac{1}{\sqrt{t}}$  and is given by  $C(0,t)$

When  $x = 2\sqrt{Dt}$ , the surface concentration has fallen by  $\frac{1}{e}$



## 7.2.5 Gaussian Solution near a Surface

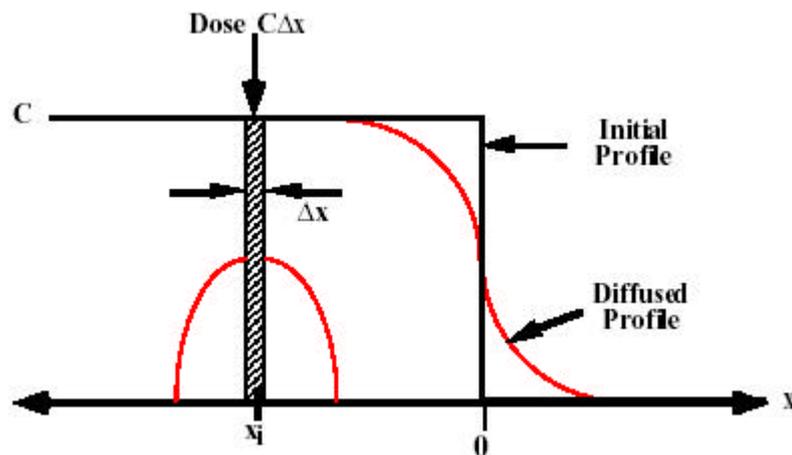


**Assumption : no dopant is lost through evaporation or segregation during the anneal**  
**Effectively a dose of 2Q is introduced into an infinite medium**

$$C(x,t) = \frac{Q}{\sqrt{pDt}} \exp\left[-\frac{x^2}{4Dt}\right] = C(0,t) \exp\left[-\frac{x^2}{4Dt}\right]$$

$$C(0,t) = \frac{Q}{\sqrt{pDt}}$$

## 7.2.6 Error-Function Solution in an Infinite Medium



$$C = 0 \text{ at } t = 0 \text{ for } x > 0$$

$$C = C \text{ at } t = 0 \text{ for } x < 0$$

The problem is made up by a sum of the previous Gaussian solutions

$$C(x,t) = \frac{C}{2\sqrt{pDt}} \sum_{i=1}^n D_{x_i} \exp - \frac{(x - x_i)^2}{4Dt}$$

$$C(x,t) = \frac{C}{2\sqrt{pDt}} \int_0^x \exp - \frac{(x - a)^2}{4Dt} da$$

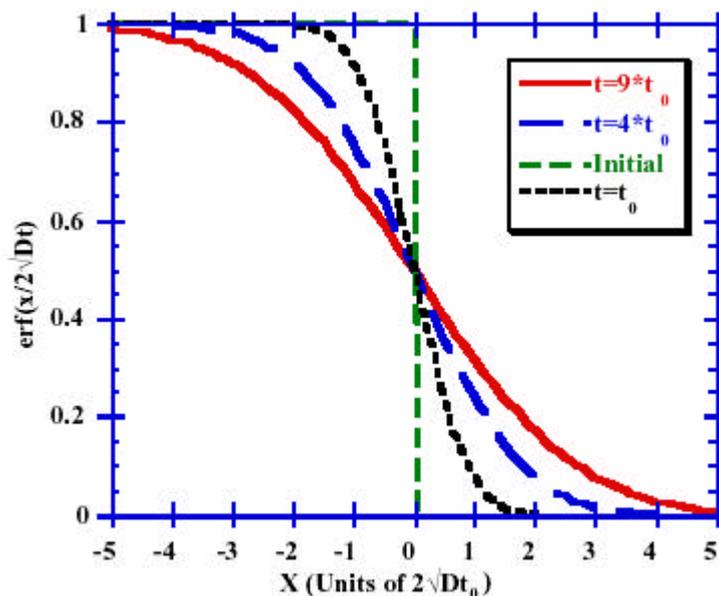
$$\frac{x - a}{2\sqrt{Dt}} = h$$

$$C(x,t) = \frac{C}{\sqrt{p}} \int_0^x \exp(-h^2) dh$$

$$\text{erf}(z) = \frac{2}{\sqrt{p}} \int_0^z \exp(-h^2) dh$$

$$C(x,t) = \frac{C}{2} \left[ 1 - \text{erf} \left( \frac{x}{2\sqrt{Dt}} \right) \right]$$

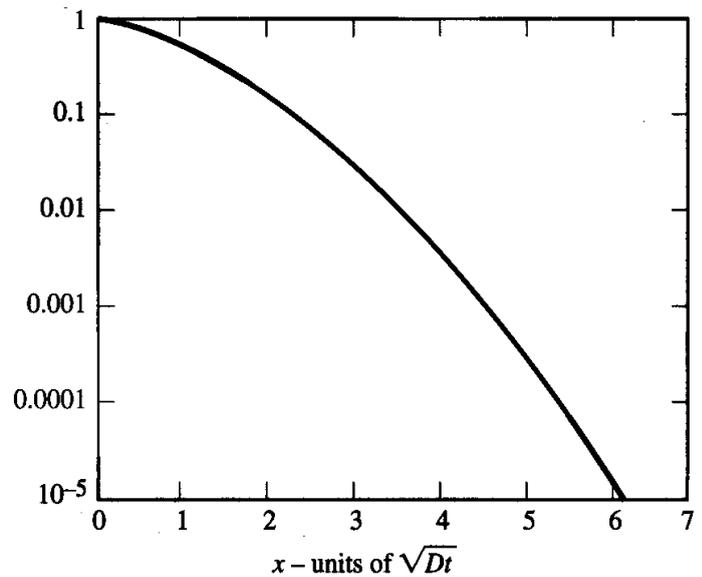
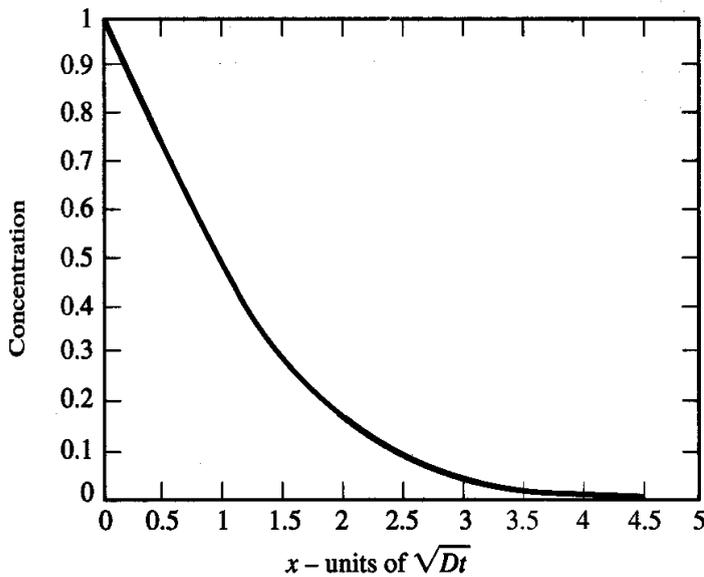
$$= \frac{C}{2} \text{erfc} \left( \frac{x}{2\sqrt{Dt}} \right)$$



### 7.2.7 Error-Function Solution near a Surface

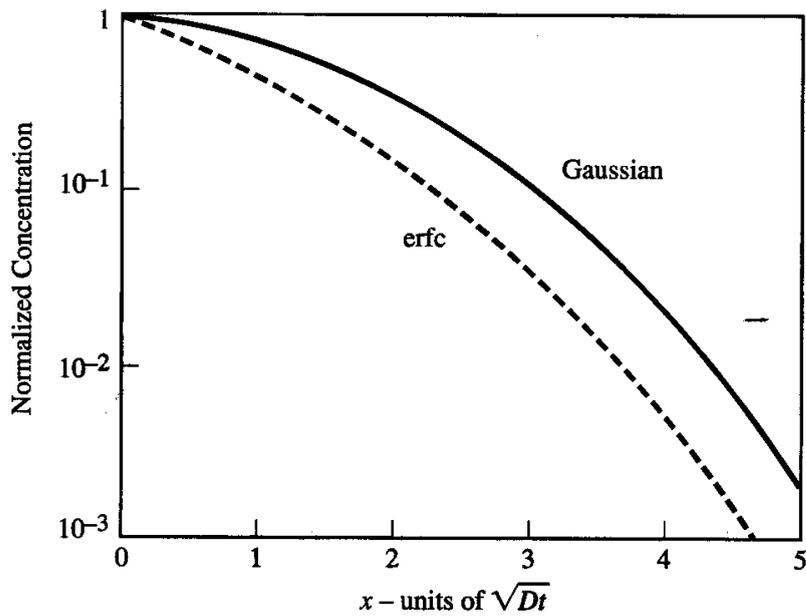
- Constant surface concentration at all times
- The midpoint in Fig.7-12 must remain stationary

$$C(x,t) = C_s \operatorname{erfc} \frac{x}{2\sqrt{Dt}}$$



- The error-function solution is (very approximately) triangular on a linear scale  $\textcircled{R} Q = C_s \sqrt{Dt}$   
More accurately,

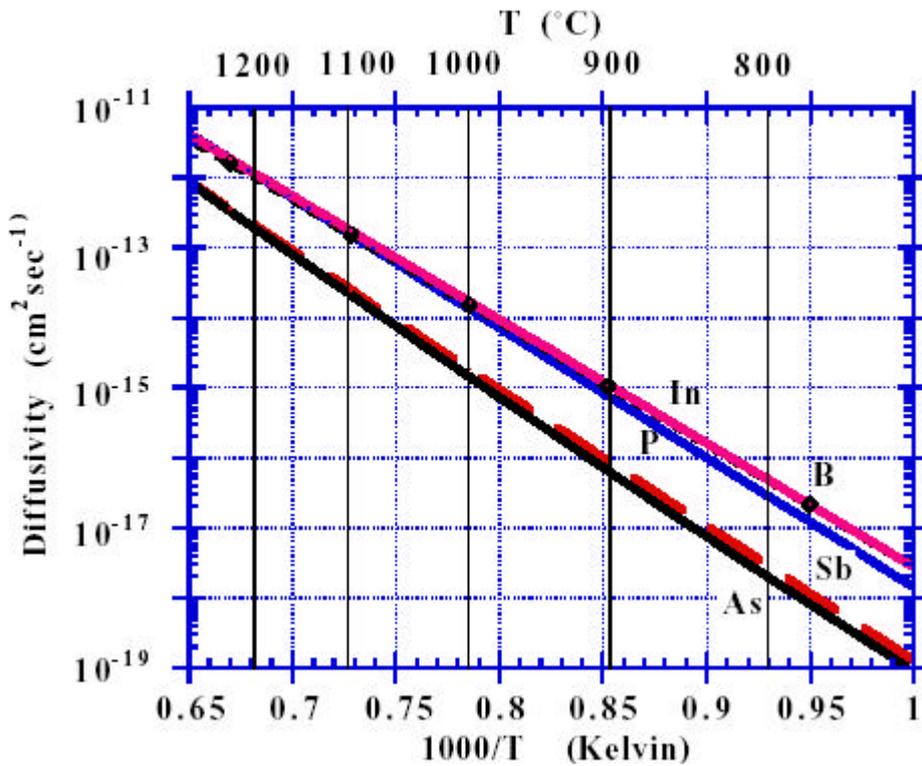
$$Q = C_s \int_0^x \operatorname{erfc} \frac{x}{2\sqrt{Dt}} dx = \frac{2C_s}{\sqrt{p}} \sqrt{Dt}$$

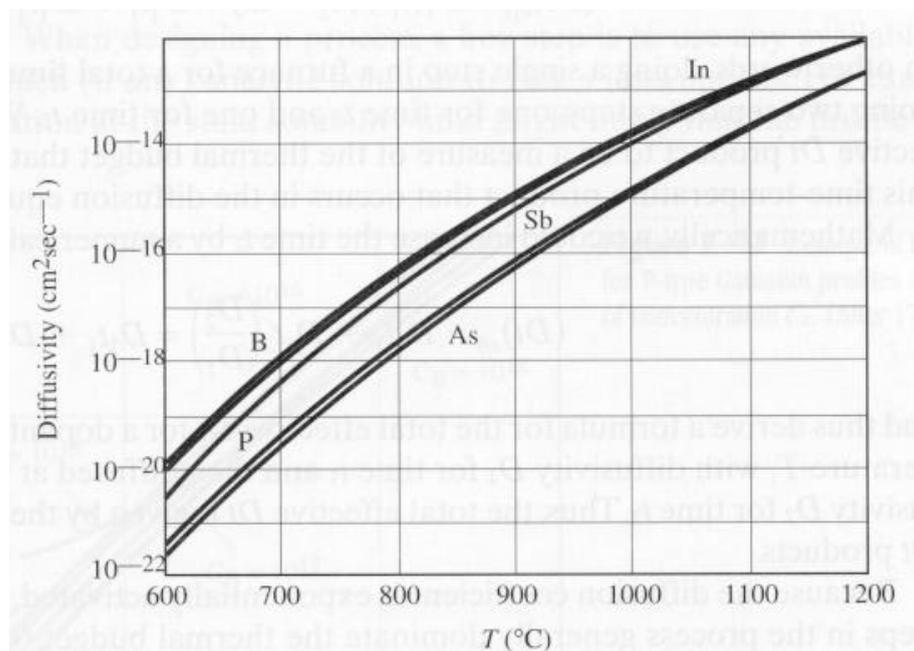


## 7.2.8 Intrinsic Diffusion Coefficients of Dopants in Silicon

$$D = D^0 \exp\left(-\frac{E_A}{kT}\right)$$

$$E_A = 3.5 - 4.5 \text{ eV}$$





**Table 7-3** Intrinsic diffusivity in  $\text{cm}^2\text{sec}^{-1}$  for silicon self-diffusion and of common dopants in single-crystal silicon, fitted to an Arrhenius expression

	Si	B	In	As	Sb	P	Units
$D^0$	560	1.0	1.2	9.17	4.58	4.70	$\text{cm}^2\text{sec}^{-1}$
$E_A$	4.76	3.5	3.5	3.99	3.88	3.68	eV

- “Slow” diffusers : As, Sb
- “Fast” diffusers : P, B, In
- Under extrinsic conditions, all dopants show much higher D

## 7.2.9 Effect of Successive Diffusion Steps

- If all the diffusion steps occurred at a constant temperature

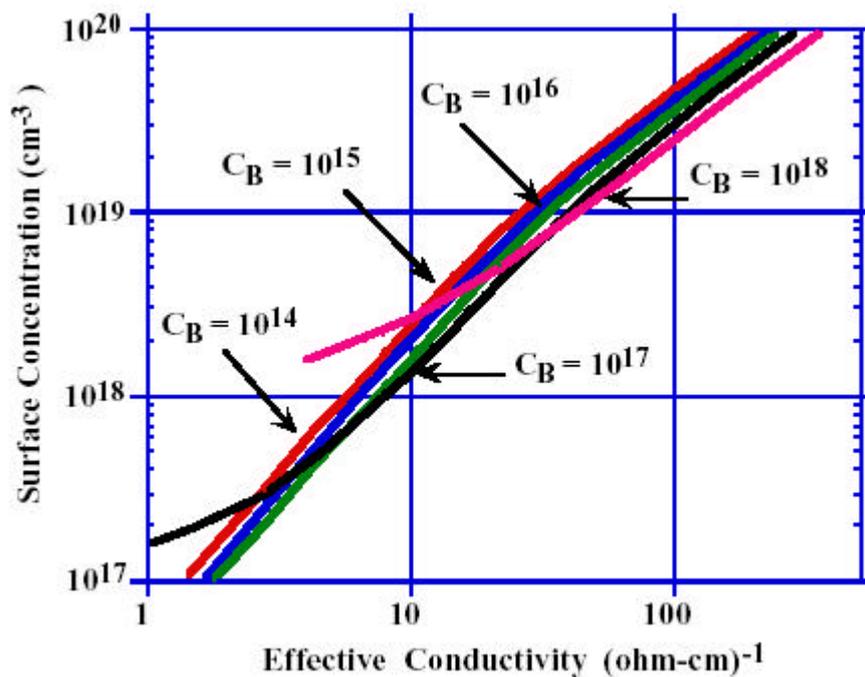
$$(Dt)_{\text{eff}} = D_1(t_1 + t_2 + \dots) = D_1t_1 + D_1t_2 + \dots$$

$$(Dt)_{\text{eff}} = D_1t_1 + D_1t_2 \frac{D_2}{D_1} = D_1t_1 + D_2t_2$$

$$t_2^{\text{equiv}} = t_1 \cdot \frac{D_1}{D_2}$$

## 7.2.10 Design and Evaluation of Diffused Layers

- Sheet resistance, surface concentration, junction depth
- Invin's curves

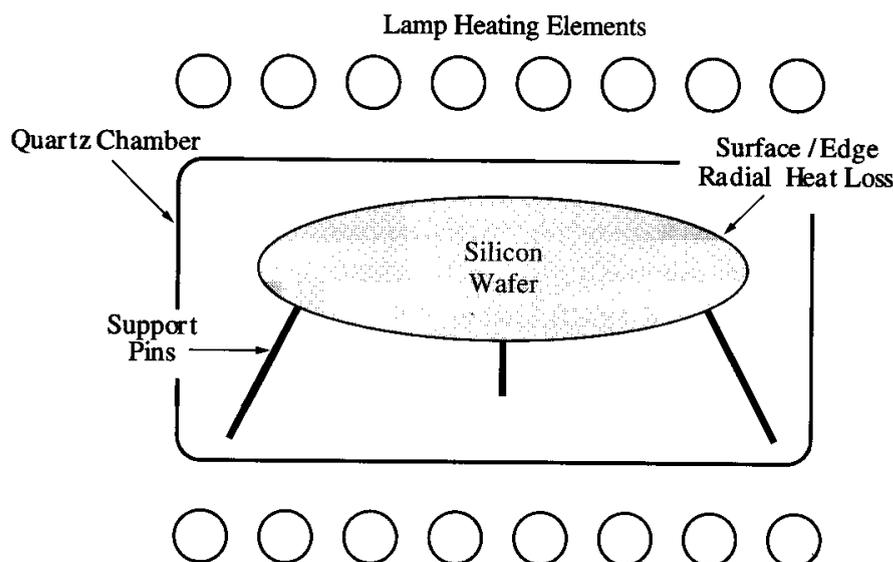


## 7.2.11 Summary of Basic Diffusion Concepts

### 7.3 Manufacturing Methods and Equipment

- Important to have some oxide or other capping layer to prevent evaporation of the dopant
- Anneals are often performed in a low partial pressure of oxygen sufficient to grow a thin oxide layer or are performed with a reasonable thickness of oxide present
- Idle at a temperature between 750-800°C
- Ramps slowly ( $5-10^{\circ}\text{C min}^{-1}$ ) to the final anneal T
- After a diffusion step, ramp down to the idling T
- Oxidation can alter D through OED
- If the dopants are introduced by ion implantation, they can end up diffusing more at lower T because the ion-implantation-induced defects last longer and continue to enhance D
- Modern devices often use a thermal cycle simply to electrically activate implanted dopant atoms and process designers would often be very happy if no diffusion took place during this step

? RTA (Rapid Thermal Annealing)



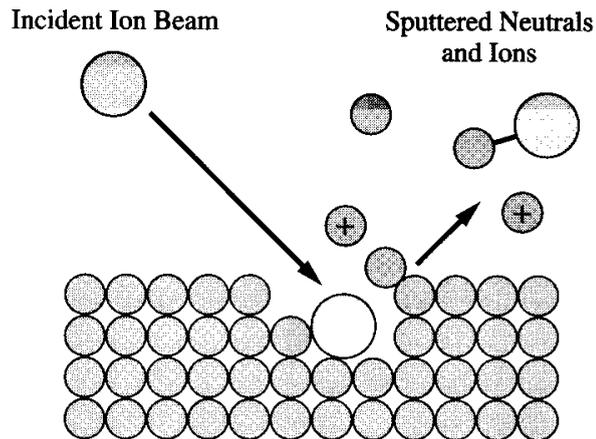
- Ramp rates ~ 100°C/sec

<b>Furnace</b>	<b>RTA</b>
Batch	Single wafer
Long time	Short time
Slow ramp	Fast ramp
Long stabilization	Short steady state
Excellent temperature control	Accurate temperature control difficult

- A silicon wafer subject to incident radiation will achieve a uniform temperature across and through the wafer in a few msecs
- Even if there is a perfectly uniform incident radiation flux, the wafer edges will be color ( heat loss)
  - ? may cause slip
  - suppressed by surrounding the wafer a polysilicon slip ring
  - multiple circular zones of lamps
- Monitoring temperature
  - Thermocouple
  - Optical pyrometry

## 7.4 Measurement Methods

### 7.4.1 SIMS (= Secondary Ion Mass Spectrometry)

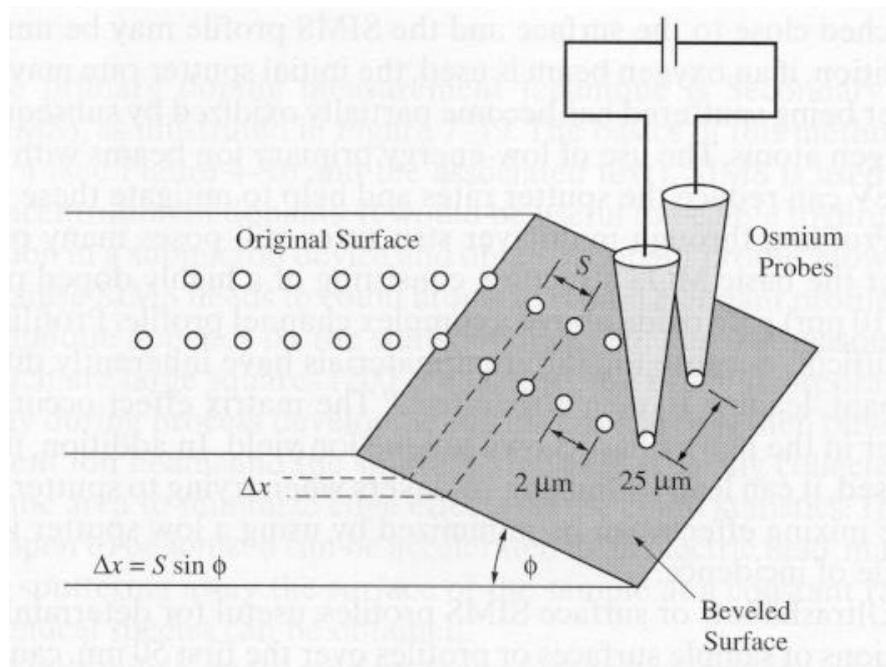


- **Special SIMS test areas, which are large squares(200×200μm), are often incorporated**
- **This large area is physically sputtered, and the ions are collected from a smaller, central probe area to minimize edge effects**
- **Accelerated ? mass analyzed ? counted**
- **The incident ion beam is chosen to produce a high yield and a small mass interference due to unintentional ion complexing**
- **n - type dopants : C<sub>s</sub>ü high yield**
- **p - type dopants : O<sub>p</sub>ÿ**
- **Sensitivity : 10<sup>16</sup>-10<sup>17</sup>/cm<sup>3</sup>**

- [ Cs : heavy ? “knock on” ? degrade the resolution  
 O : light ? minimize “knock on” problem  
 but, a higher background floor
- May be unreliable in ultrashallow
- O : the initial rate may be higher  
 ? low energy (200eV to 5 keV)
- “Matrix effect” ( each different layer in the matrix has  
 its own unique ion yield)
- Cs : can lead to “mixing” of layers when pattering through  
 ultrathin layers
- Mixing effects can be minimized by using a low energy and  
 shallow angle of incidence
- Sputter neutral mass spectrometry

#### 7.4.2 Spreading Resistance

- Electrical measurement using a pair of fine metal probes
- By measuring the resistivity and comparing it against  
 resistivity standards ?  $N$  can be obtained



- **Beveling** : gluing a small sample to a precision-machined holder  
 ? polished using a diamond slurry and a rotating wheel
- **Important** : the probes make good electrical contact yet not penetrate too deeply into the bulk
- **SIMS** : chemical  
 Spreading resistance : electrical } difference corresponds to dopants which are not electrically active

### 7.4.3 Sheet Resistance

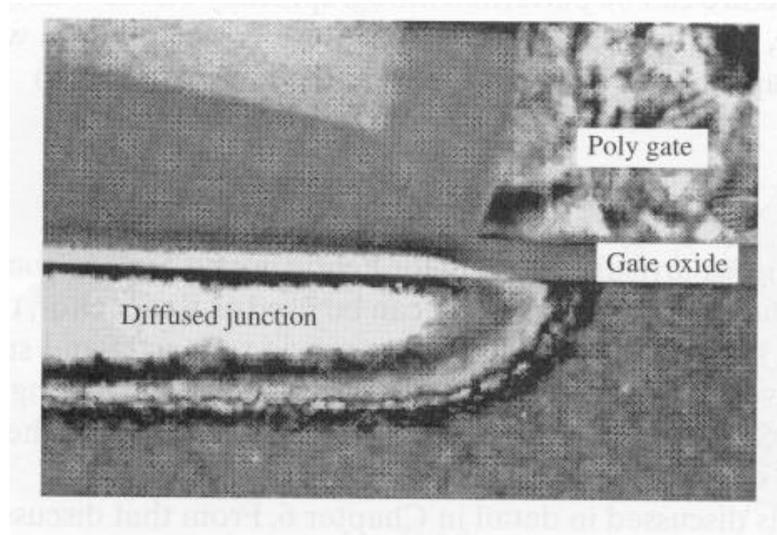
- Provide an integrated electrical conductivity over the whole profile
- 4 probes  $\left[ \begin{array}{l} \text{outer 2 : force a current} \\ \text{inner 2 : measure a voltage} \end{array} \right.$
- For shallow junctions, the mechanical probes penetrate too deeply  
 ? Van der Pauw

### 7.4.4 Capacitance Voltage

$C$  ?  $x_D$  ?  $N$

### 7.4.5 TEM Cross Section

- TEM, electrical inverse methods, direct CV, Kelvin probe



- HF : HNO<sub>3</sub> : CH<sub>3</sub>COOH (1 : 40 : 20)  
↓                      ↘                      ↗  
                                 oxidizes the Si                      dilutant

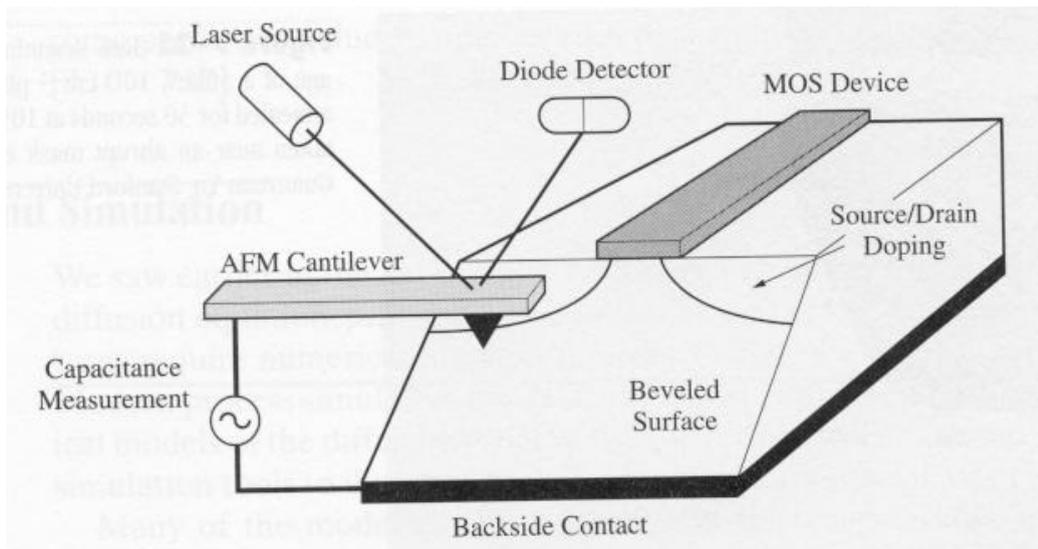
dissolves the oxide

- The etch rates depend on N and appear to similar for n-type and p-type

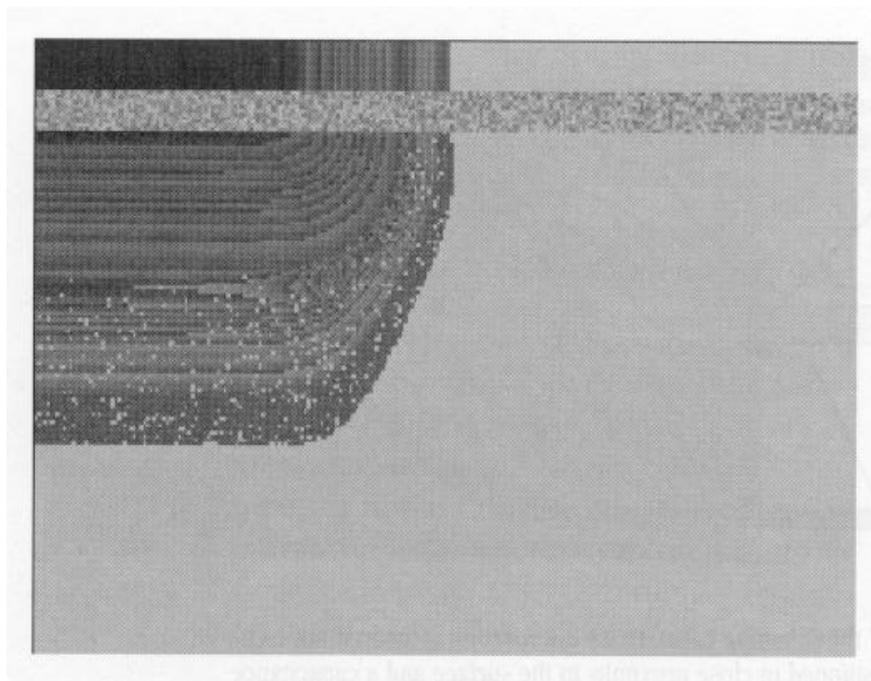
### 7.4.6 2D Electrical Measurements Using Scanning Probe Microscopy

- Average distance  $\frac{\lambda}{1} = 1.3\text{nm}$  at  $10^{20}\text{cm}^{-3}$   
 $\frac{\lambda}{1} = 6.2\text{nm}$  at  $10^{18}\text{cm}^{-3}$   
 $\frac{\lambda}{1} = 28.8\text{nm}$  at  $10^{16}\text{cm}^{-3}$

- **Scanning probe method ? STM (Scanning Tunneling Microscope)**
- **Scanning capacitance and scanning resistance probes**



- **Cross-sectional samples must be prepared**
- **The surface of the Si cross section or metal probe must be coated with an oxide layer**



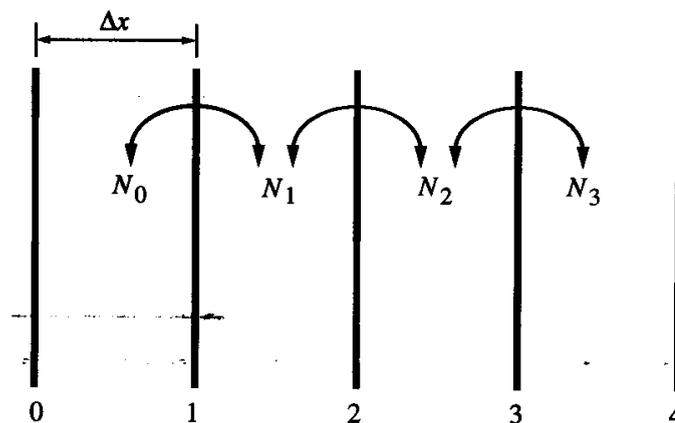
- SRP
- Workfunction measurement ? N

### 7.4.7 Inverse Electrical Measurements

- IV or CV  $\longleftrightarrow$  simulation results compared
- The doping profiles are varied until they match

## 7.5 Models and Simulation

### 7.5.1 Numerical Solutions of the Diffusion Equation



$$C_i = N_i \cdot x / \text{cm}^3$$

- The atoms remain relatively fixed in the lattice but vibrate at the Debye frequency ( $10^{13} \text{sec}^{-1}$ )

- **The frequency of hopping**

$$v_b = v_d \exp\left\{-\frac{E_b}{kT}\right\}$$

↓  
**Debye frequency**

$$\frac{v_b}{2} N \text{ jumping right}$$

$$\frac{v_b}{2} N \text{ jumping left}$$

$$F = -\frac{v_b}{2} (N_2 - N_1) = -\frac{v_b}{2} D_x (C_2 - C_1) = -\frac{v_b}{2} D_x^2 \frac{DC}{Dx}$$

$$= -D \frac{DC}{Dx}$$

Where  $D = \frac{v_b}{2} D_x^2$

- **At plain i**

$$N_i^+ = N_i + \frac{v_b}{2} Dt (N_{i-1} - 2N_i + N_{i+1})$$

$$C_i^+ = C_i + \frac{v_b}{2} Dt (C_{i-1} - 2C_i + C_{i+1})$$

$$C_i^+ = C_i + \frac{DDt}{Dt^2} (C_{i-1} - 2C_i + C_{i+1})$$

- **A simplification of this equation ;**

$$\frac{DDt}{Dx^2} = \frac{1}{2}$$

$$C_i^+ = \frac{1}{2} (C_{i-1} + C_{i+1})$$

- If  $\frac{DDt}{Dx^2} > 0.5$  @ the solution becomes unstable
- $\frac{DDt}{Dx^2} = 0.5$  represents the maximum value

### 7.5.2 Modification to Fick's Laws to Account for Electric Field Effects

- When  $N > n_i$ , e-fields set up by the dopants can affect the diffusion process

? Additional flux or flow of species  $F' = Cv$

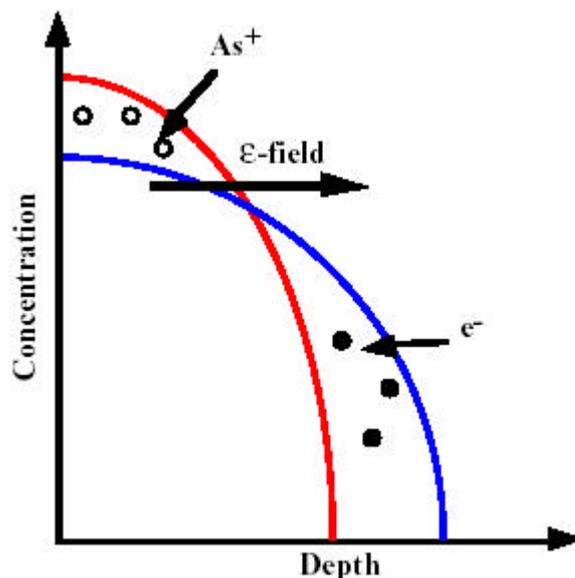
↙  
velocity of particles

$$F_{\text{total}} = F + F' = -D \frac{\partial C}{\partial x} + Cv$$

- If  $F'$  and  $v$  are independent of  $x$ ,

$$\frac{\partial C}{\partial t} = \frac{\partial}{\partial x} \left( D \frac{\partial C}{\partial x} \right) - v \frac{\partial C}{\partial x}$$

- The origin of the field comes from the higher mobility of electrons and holes compared to dopants



$$v = me$$

$$e = \frac{J_y}{J_x}$$

$$y = - \frac{kT}{q} \ln \frac{n}{n_i}$$

$$m = \frac{q}{kT} D$$

$$\begin{aligned} F &= - D \frac{J_C}{J_x} - DC \frac{J}{J_x} \ln \frac{n}{n_i} \\ &= - DC \frac{J}{J_x} \ln \frac{n}{n_i} \end{aligned}$$

- The flux can be expressed in terms of N only, if charge neutrality and Boltzmann statistics apply

$$N_D^+ + P = N_A^- + n$$

$$np = n_i^2$$

If  $N_D^+ \approx N_D$  and  $N_A^- \approx N_A$  the flux can be written as

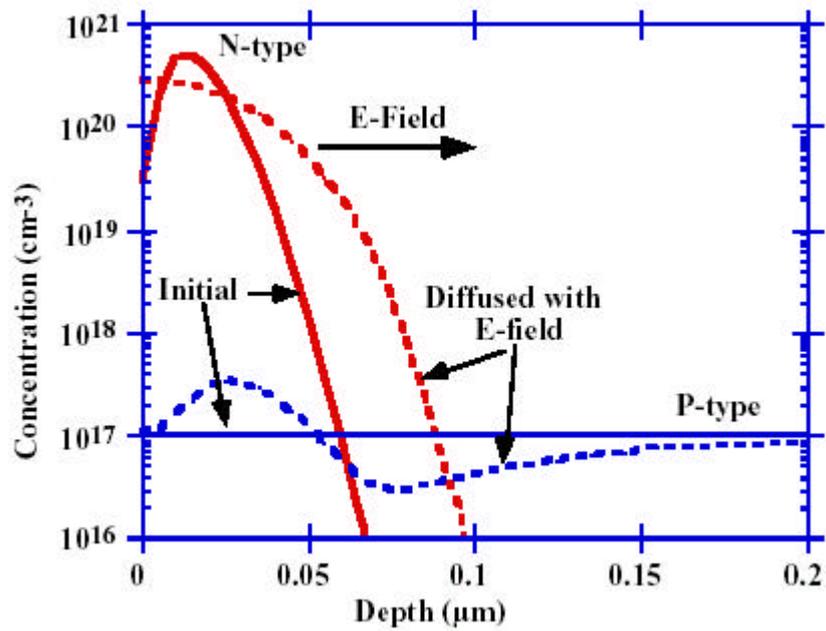
$$F = - hD \frac{J_C}{J_x}$$

where  $C = |N_D - N_A|$

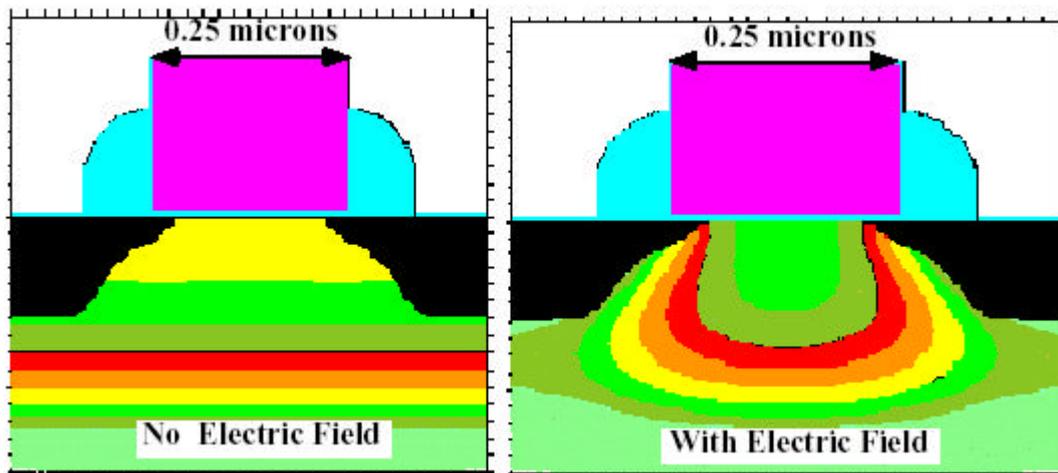
$h$  = electric field enhancement factor

$$h = 1 + \frac{C}{\sqrt{C^2 + 4n_i^2}}$$

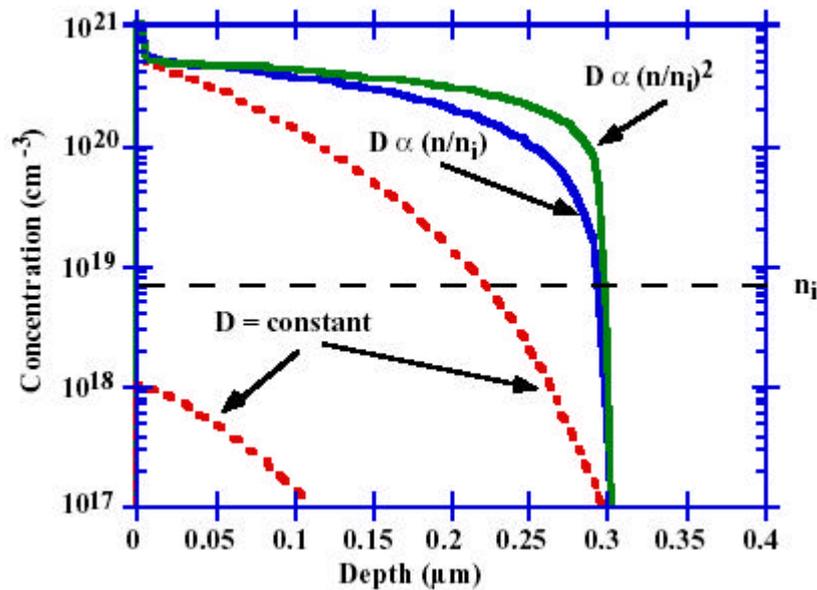
$h$  has an upper bound of 2



- Bigger changes in D of a low-concentration dopant



### 7.5.3 Modifications to Fick's Laws to Account for Concentration-Dependent Diffusion



- The diffusivity varies with concentration, with the higher-concentration regions diffusing faster

$$\frac{\partial C}{\partial t} = \frac{\partial}{\partial x} \left( D_{\text{eff}} \frac{\partial C}{\partial x} \right)$$

- Experimentally,  $D \propto \frac{n}{n_i}$  or  $D \propto \left(\frac{n}{n_i}\right)^2$

$$D_{\text{eff}} = D^0 + D^- \frac{n}{n_i} + D^+ \left(\frac{n}{n_i}\right)^2 \quad \text{for Ntype dopants}$$

$$D_{\text{eff}} = D^0 + D^+ \frac{P}{n_i} + D^{++} \left(\frac{P}{n_i}\right)^2 \quad \text{for Ptype dopants}$$

- Intrinsic Conditions**

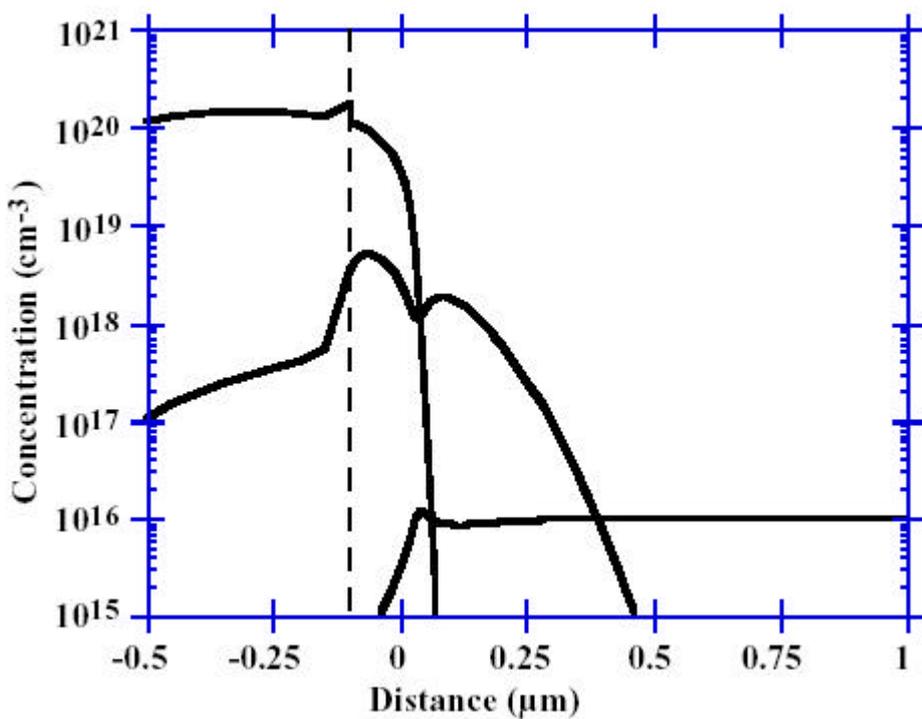
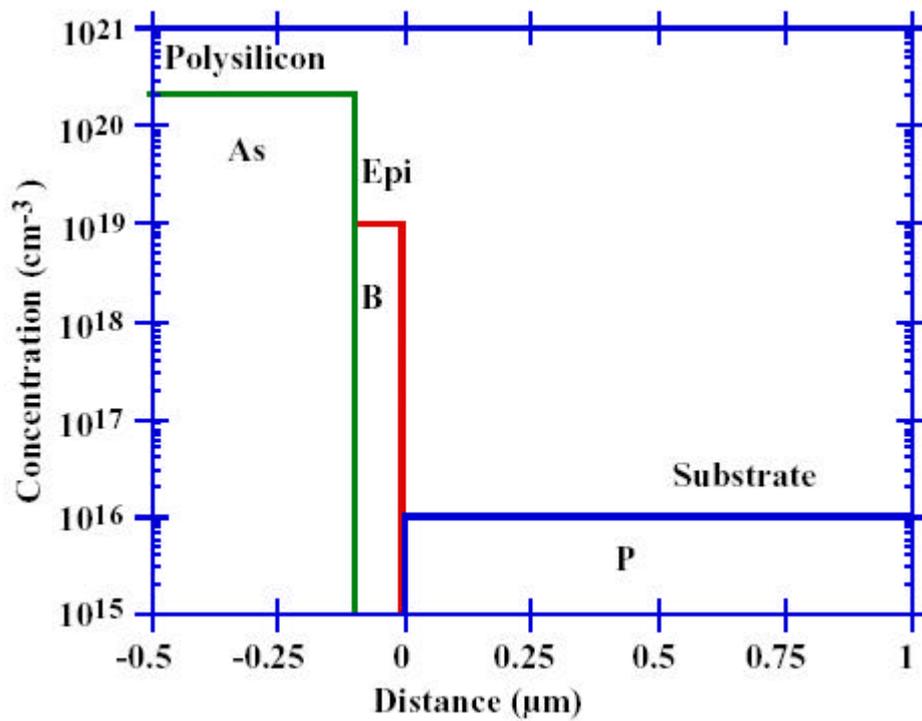
$$D_A^* = D^0 + D^- + D^=$$

$$D = D^0 \exp\left\{-\frac{E_D}{kT}\right\}$$

$$D_A^{\text{eff}} = D_A^* \frac{1 + b \frac{n}{n_i} + g \frac{n^2}{n_i^2}}{1 + b + g}$$

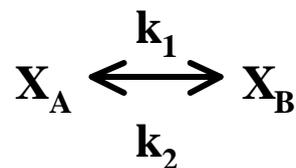
where  $b = \frac{D^-}{D^0}$  and  $g = \frac{D^=}{D^0}$

	Si	B	In	As	Sb	P
$D^0$ .0 cm <sup>2</sup> sec <sup>-1</sup>	560	0.05	0.6	0.011	0.214	3.85
$D^0$ .E eV	4.76	3.5	3.5	3.44	3.65	3.66
$D^+$ .0 cm <sup>2</sup> sec <sup>-1</sup>		0.95	0.6			
$D^+$ .E eV		3.5	3.5			
$D^-$ .0 cm <sup>2</sup> sec <sup>-1</sup>				31.0	15.0	4.44
$D^-$ .E eV				4.15	4.08	4.0
$D^=$ .0 cm <sup>2</sup> sec <sup>-1</sup>						44.2
$D^=$ .E eV						4.37



## 7.5.4 Segregation

- Dopants have different solubilities in different materials  
? redistribute at an interface until the chemical potential is the same
- Segregation coefficient = ratio of the equilibrium doping concentration on each side of the interface

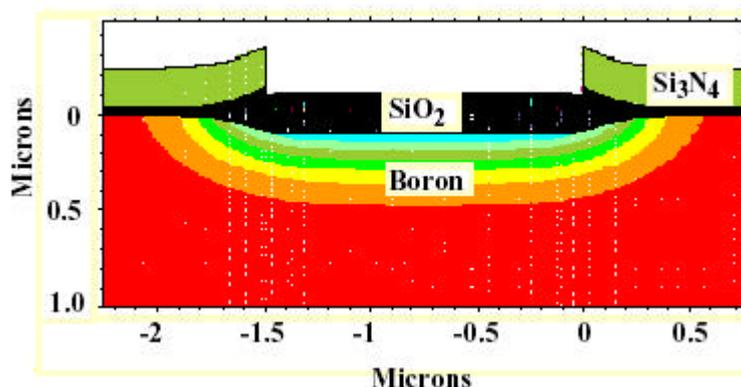


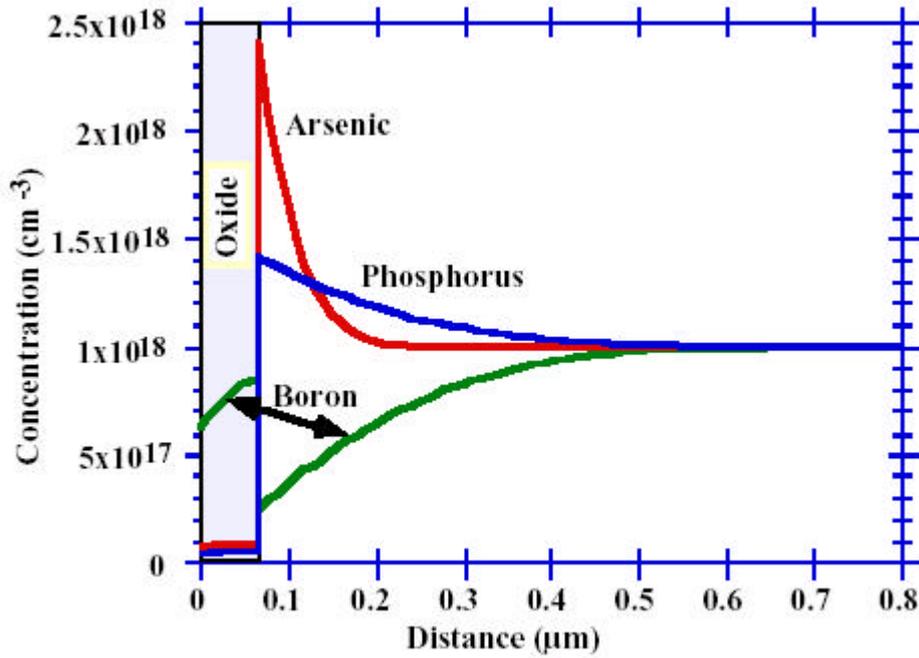
$$F_{AB} = k_1 C_A - k_2 C_B$$

In steady state,  $F_{AB} = 0 \Rightarrow \frac{C_B}{C_A} = \frac{k_1}{k_2} = k_0 = \text{segregation coefficient}$

$$\nabla F = k_1 \frac{\partial C_A}{\partial x} - \frac{C_B}{k_0} \frac{\partial}{\partial x} = k_1 \frac{\partial C_A}{\partial x} - \frac{C_B}{k_0} \frac{\partial}{\partial x}$$

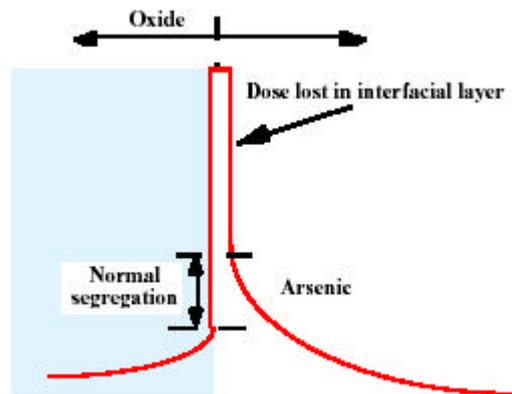
$$k_0 = \frac{C_{\text{Si}}}{C_{\text{SiO}_2}} \gg \begin{cases} 1 & \text{for B} \\ 10 & \text{for As, Sb, P} \end{cases}$$



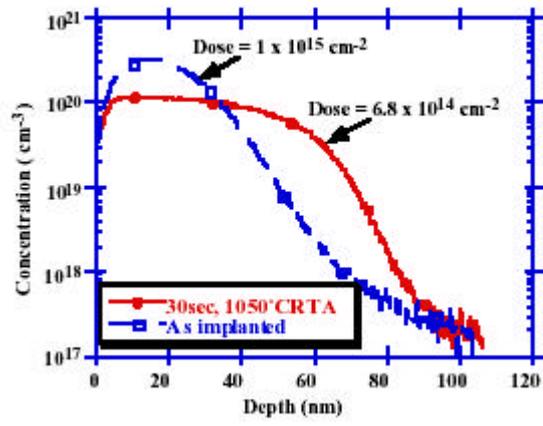


- As profile is steeper than the P profile because of smaller diffusivity

### 7.5.5 Interfacial Dopant Pileup



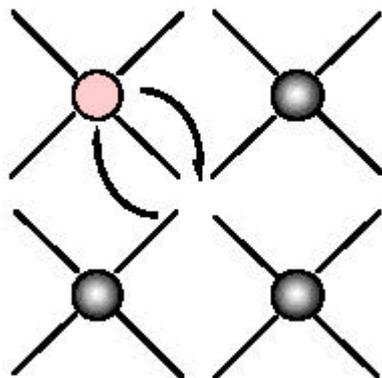
- The interfacial layer may be as thin as a monolayer and can act as a sink for dopant atoms and is able to trap on the order of an atomic layer of dopant



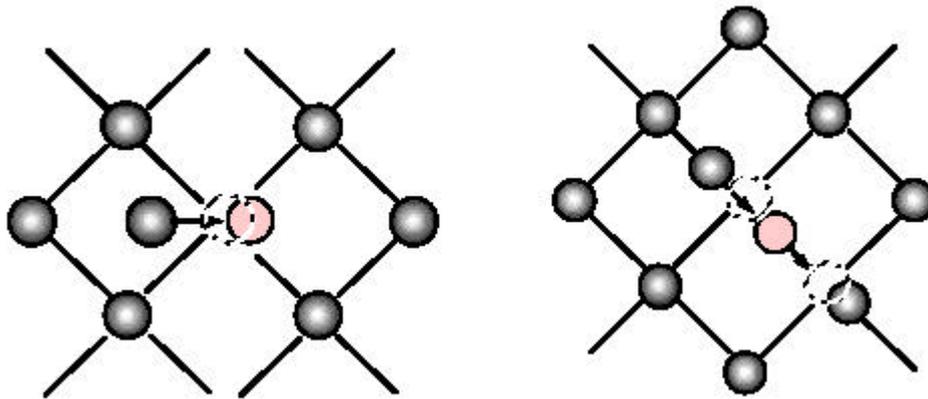
- It occurs on the same time scale as the rapid thermal cycles

### 7.5.6 Summary of the Macroscopic Diffusion Approach

### 7.5.7 The Physical Basic for Diffusion at an Atomic Scale



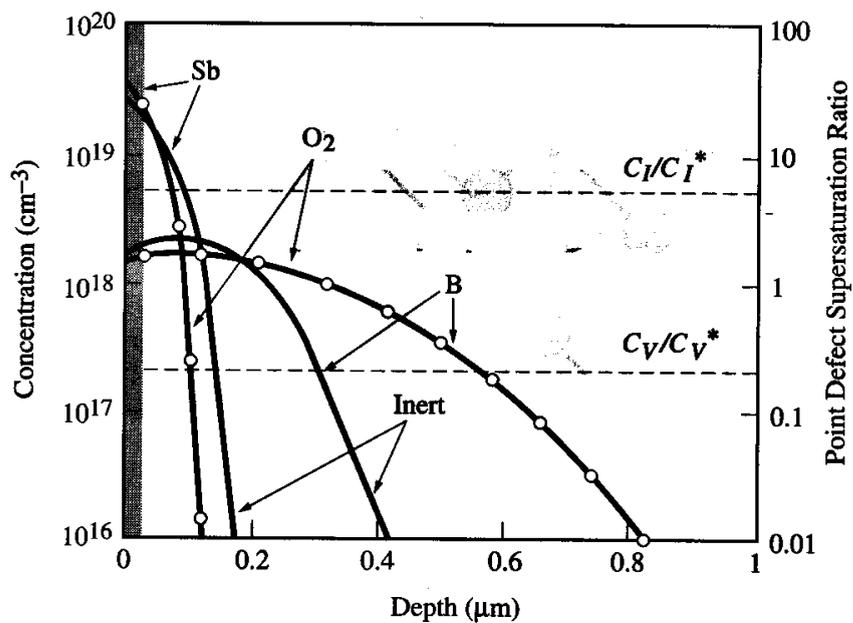
Vacancy-assisted diffusion



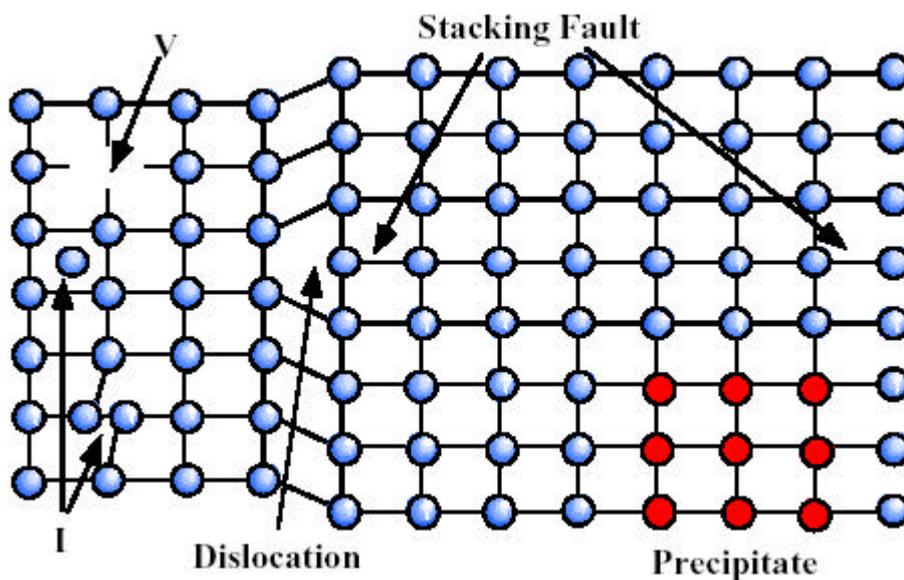
Interstitial-assisted } kick-out  
interstitialcy

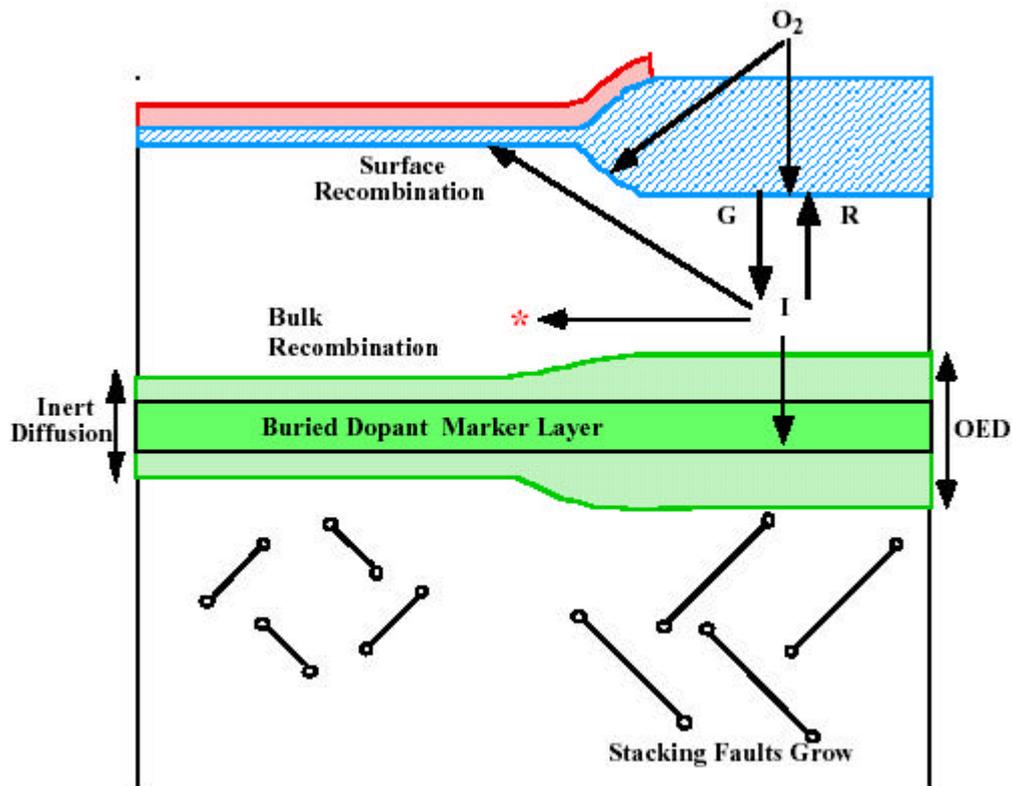
### 7.5.8 Oxidation-Enhanced or –Retarded Diffusion

- P, B : Oxidation-enhanced diffusion (OED)
- Sb : Oxidation-retarded diffusion (ORD)

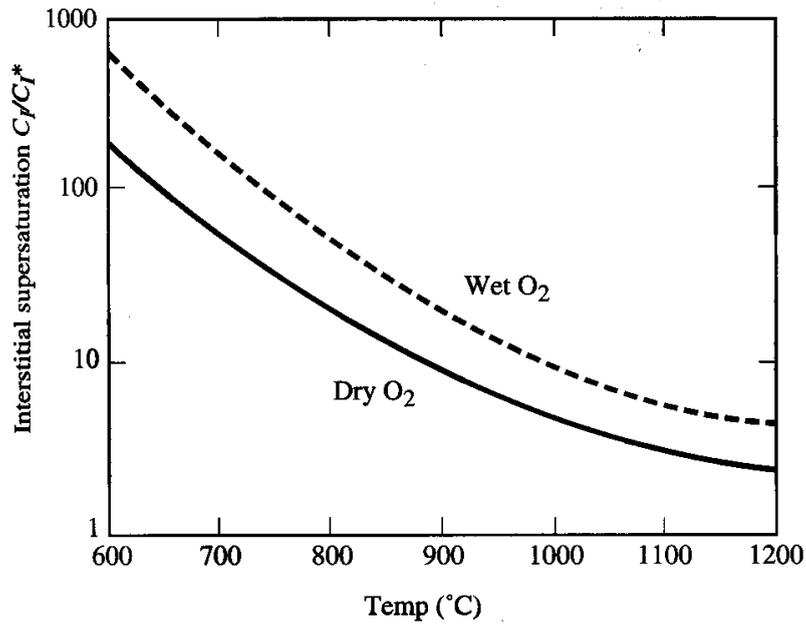


- **Perturbed point-defect concentrations caused by the surface oxidation**
- **Growth of oxidation stacking faults**  $\frac{\dot{u}}{\dot{y}}$  same underlying phenomena  
OED  $\dot{p}$
- **Stacking faults = nonequilibrium defect structures composed of an extra partial plane of interstitials**





- When Si is oxidized ? expansion of 30% ? compressive stress ? injection of interstitial to make space at the Si surface
  - ↙ large size
- Sb prefers to diffuse with vacancies. The interstitials recombine with vacancies ? vacancy concentration?
  - ? D?

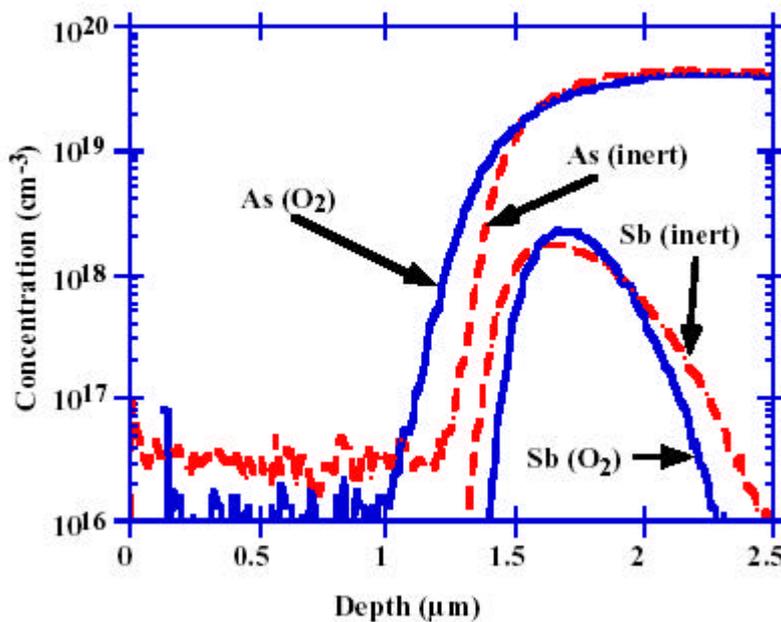


- The interstitial level depends on the balance between G and R

$$C_i \propto \sqrt{\frac{dx}{dt}}$$

- Interstitial supersaturation is larger at lower oxidation T ? enhancement in D

### 7.5.9 Dopant Diffusion Occurs by Both I and V



- The diffusion of  $\begin{cases} \text{As is enhanced} \\ \text{Sb is retarded} \end{cases}$
- Nitridation of the silicon surface in an ammonia has exactly the opposite effects
  - $\begin{cases} \text{B \& P : retarded diffusion} \\ \text{Sb : enhanced} \end{cases}$
- Nitridation injects vacancies
- Dopants diffuse with a fraction  $f_I$  interstitial-type diffusion and with a fraction  $f_V=1-f_I$  vacancy-type mechanism

$$D_A^{\text{eff}} = D_A^* \frac{C_I}{C_I^*} f_I + f_V \frac{C_V}{C_V^*} \ddot{\theta}$$

$$\begin{aligned} \frac{D}{D^*} \ddot{\theta}^R &= f_I^R \frac{C_I}{C_I^*} + f_V^R \frac{C_V}{C_V^*} \\ &= (1 - f_V^R) \frac{C_I}{C_I^*} + f_V^R \frac{C_V}{C_V^*} \end{aligned}$$

- A very conservative assumption :  $\frac{D}{D^*} \ddot{\theta}^R \approx (1 - f_V^R) \frac{C_I}{C_I^*}$

An even more conservative assumption :  $f_V^R > 1 - \frac{D}{D^*} \ddot{\theta}^R$

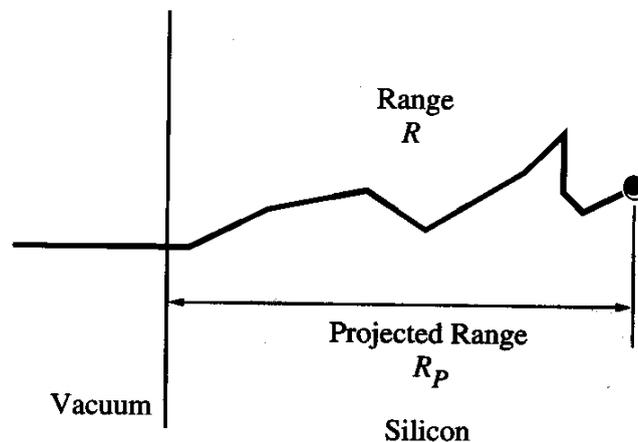
## Chap 8 Ion Implantation

### 8.1 Introduction

- Create a cascade of damage that may displace a thousand silicon atoms for each implanted ion

### 8.2 Historical Development and Basic Concepts

- Provides a very precise means to introduce a specific dose ( the electrical charge on the ion allows them to be counted by collection in a Faraday cup)
- I/I is a random process



- Large numbers of ions are implanted so an average depth for the implanted dopants can be calculated
- Heavy ions do not travel as far as light ions